

Lagrangian Descriptors: A powerful method for investigating the behavior and chaoticity of dynamical systems

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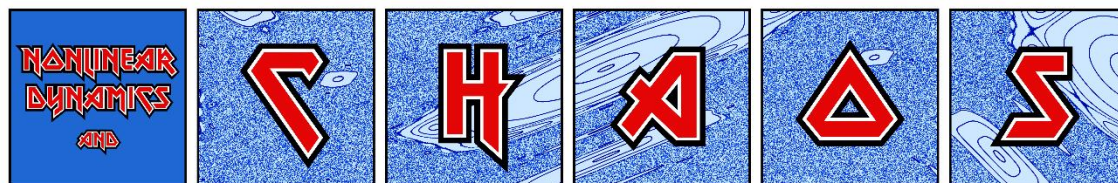
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Outline

- Lagrangian descriptors (LDs)
- Smaller Alignment Index (SALI)
- Chaos diagnostics based on LDs:
 - ✓ the difference of LDs of neighboring orbits
 - ✓ the ratio of LDs of neighboring orbits
- Applications:
 - ✓ Hénon – Heiles system
 - ✓ 2D Standard map
 - ✓ Motion of a satellite
- Summary

Lagrangian descriptors (LDs)

The computation of LDs is based on the accumulation of some positive scalar value along the path of individual orbits.

Consider an N dimensional continuous time dynamical system

$$\dot{x} = \frac{dx(t)}{dt} = f(x, t)$$

The arclength definition (Madrid, Mancho, Chaos, 2009 – Mendoza, Mancho, PRL, 2010 – Mancho et al., Commun. Nonlin. Sci. Num. Simul., 2013).

Forward time LD :

$$LD^f(x, \tau) = \int_0^\tau \|\dot{x}(t)\| dt$$

Backward time LD :

$$LD^b(x, \tau) = \int_{-\tau}^0 \|\dot{x}(t)\| dt$$

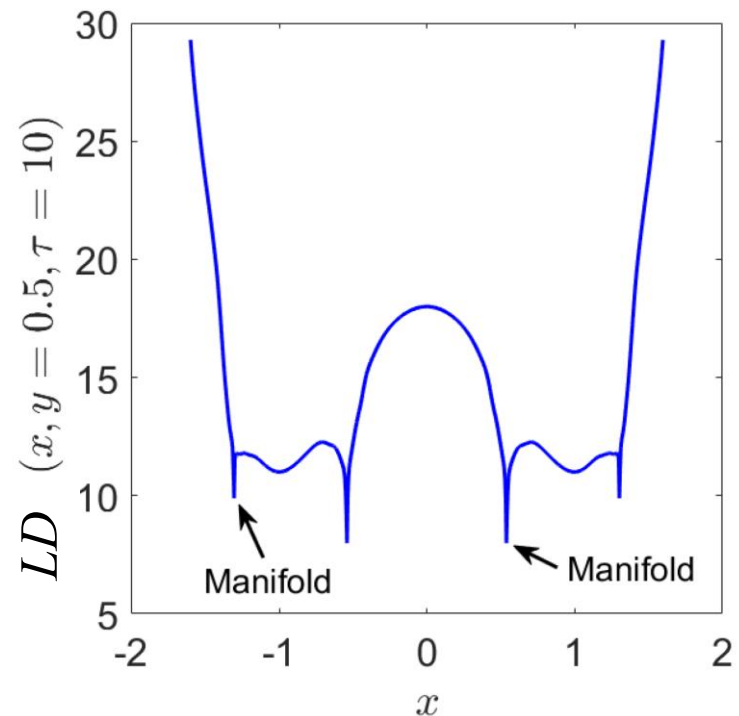
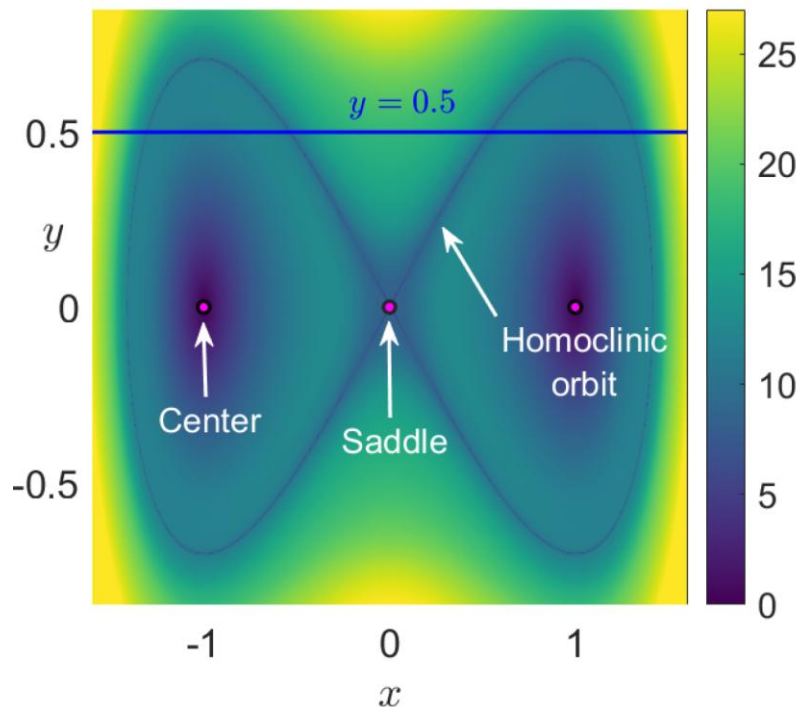
Combined LD :

$$LD(x, \tau) = LD^b(x, \tau) + LD^f(x, \tau)$$

LDs: 1 dof Duffing Oscillator

$$H(x, y) = \frac{1}{2}y^2 + \frac{1}{4}x^4 - \frac{1}{2}x^2$$

The system has three equilibrium points: a saddle located at the origin and two diametrically opposed centers at the points $(\pm 1, 0)$.



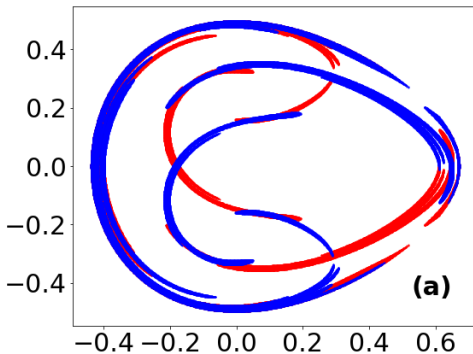
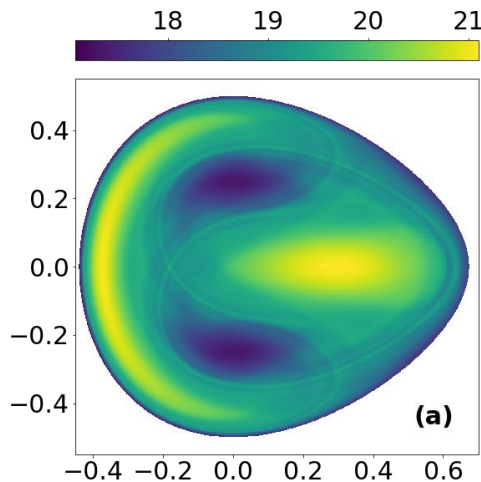
From Agaoglou et al. 'Lagrangian descriptors: Discovery and quantification of phase space structure and transport', 2020, <https://doi.org/10.5281/zenodo.3958985>

The **location of the stable and unstable manifolds** can be extracted from the ridges of the **gradient field of the LDs** since they are located at **points where the forward and the backward components of the LD are non-differentiable**.

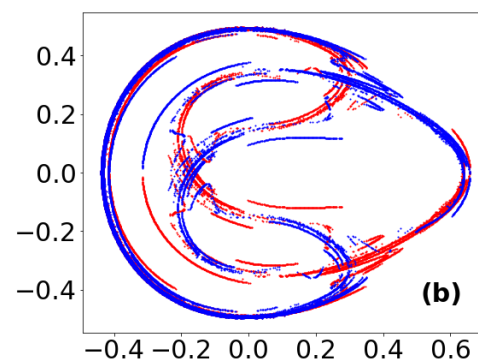
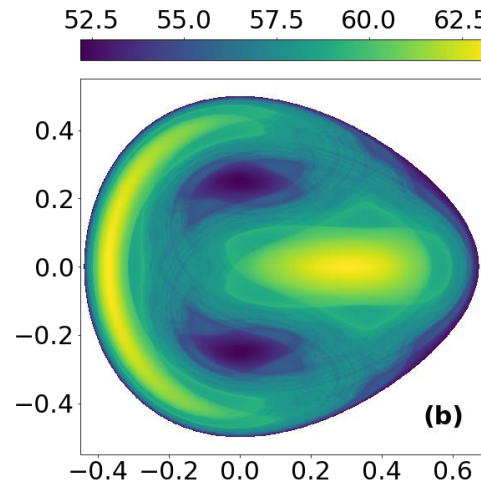
Lagrangian descriptors (LDs)

Hénon-Heiles system: $H = \frac{1}{2}(p_x^2 + p_y^2) + \frac{1}{2}(x^2 + y^2) + x^2y - \frac{1}{3}y^3$

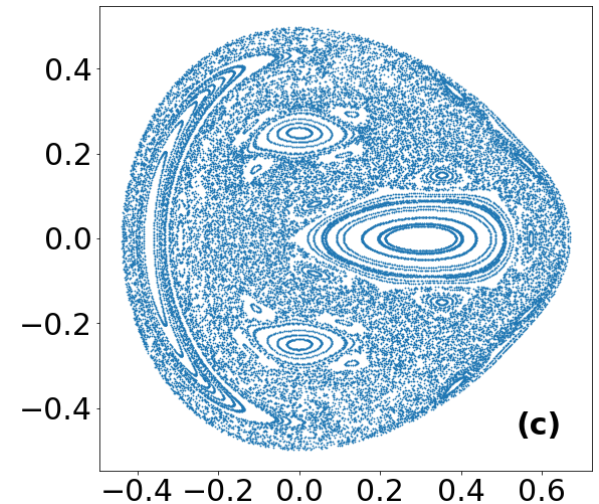
$\tau=20$



$\tau=60$



$H=1/8$



Stable and unstable manifolds

Lagrangian descriptors (LDs)

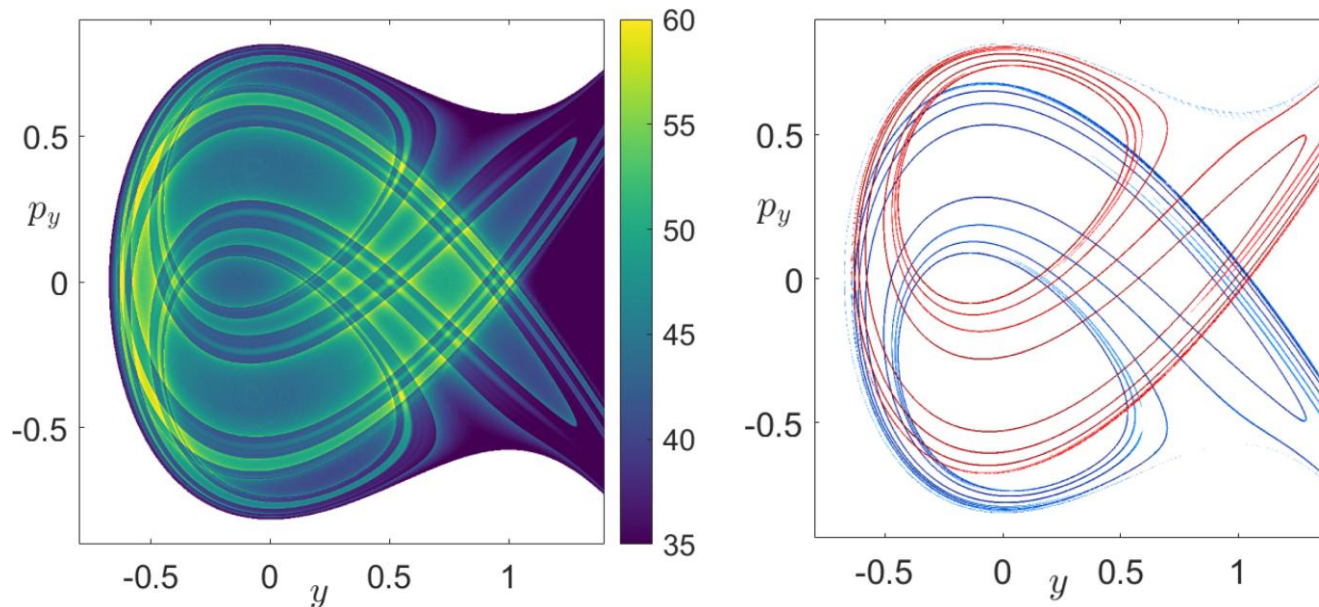
The '*p*-norm' definition (Lopesino et al., Commun. Nonlin. Sci. Num. Simul., 2015 – Lopesino et al., Int. J. Bifurcat. Chaos, 2017).

Combined *LD* (usually $p=1/2$):

$$LD(x, \tau) = \int_{-\tau}^{\tau} \left(\sum_{i=1}^N |f_i(x, t)|^p \right) dt$$

Hénon-Heiles system: $H = \frac{1}{2}(p_x^2 + p_y^2) + \frac{1}{2}(x^2 + y^2) + x^2y - \frac{1}{3}y^3$

Stable and unstable manifolds for $H=1/3$, $\tau=10$.



Maximum Lyapunov Exponent (MLE)

Chaos: sensitive dependence on initial conditions.

Roughly speaking, the MLE of a given orbit characterizes the **mean exponential rate of divergence** of trajectories surrounding it.

Consider an orbit in the $2N$ -dimensional phase space with **initial condition $x(0)$** and **an initial deviation vector (small perturbation) from it $v(0)$** .

Then the mean exponential rate of divergence is:

$$MLE = \lambda_1 = \lim_{t \rightarrow \infty} \Lambda(t) = \lim_{t \rightarrow \infty} \frac{1}{t} \ln \frac{\|v(t)\|}{\|v(0)\|}$$

$\lambda_1 = 0 \rightarrow$ Regular motion ($\Lambda \propto t^{-1}$)

$\lambda_1 > 0 \rightarrow$ Chaotic motion

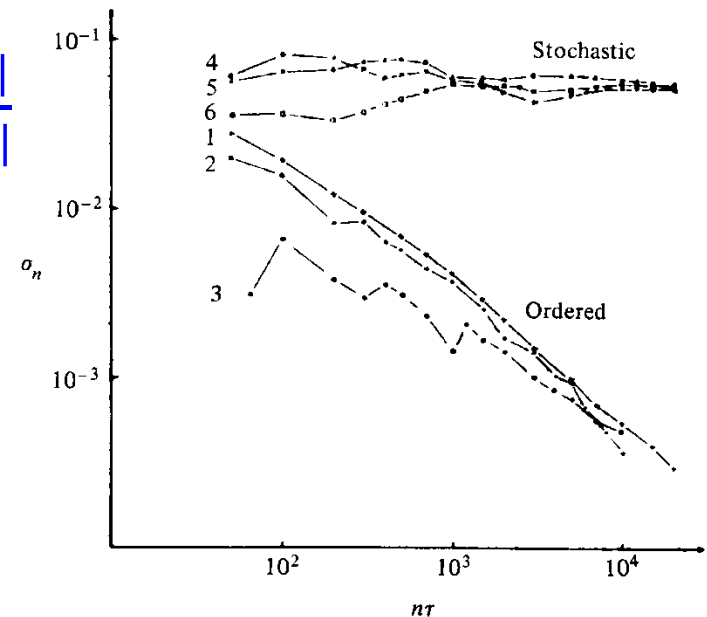


Figure 5.7. Behavior of σ_n at the intermediate energy $E = 0.125$ for initial points taken in the ordered (curves 1–3) or stochastic (curves 4–6) regions (after Benettin *et al.*, 1976).

The Smaller Alignment Index (SALI)

Consider the $2N$ -dimensional phase space of a conservative dynamical system (symplectic map or Hamiltonian flow).

An orbit in that space with initial condition :

$$P(0) = (x_1(0), x_2(0), \dots, x_{2N}(0))$$

and a deviation vector

$$v(0) = (\delta x_1(0), \delta x_2(0), \dots, \delta x_{2N}(0))$$

The evolution in time (in maps the time is discrete and is equal to the number n of the iterations) of a deviation vector is defined by:

- the variational equations (for Hamiltonian flows) and
- the equations of the tangent map (for mappings)

Definition of the SALI

We follow the evolution in time of two different initial deviation vectors ($v_1(0)$, $v_2(0)$), and define SALI [S., J. Phys. A (2001) – S. & Manos, Lect. Notes Phys. (2016)] as:

$$SALI(t) = \min\{\|\hat{v}_1(t) + \hat{v}_2(t)\|, \|\hat{v}_1(t) - \hat{v}_2(t)\|\}$$

where

$$\hat{v}_1(t) = \frac{v_1(t)}{\|v_1(t)\|}$$

When the two vectors become collinear

$$SALI(t) \rightarrow 0$$

SALI – Hénon-Heiles system

As an example, we consider the 2D Hénon-Heiles system:

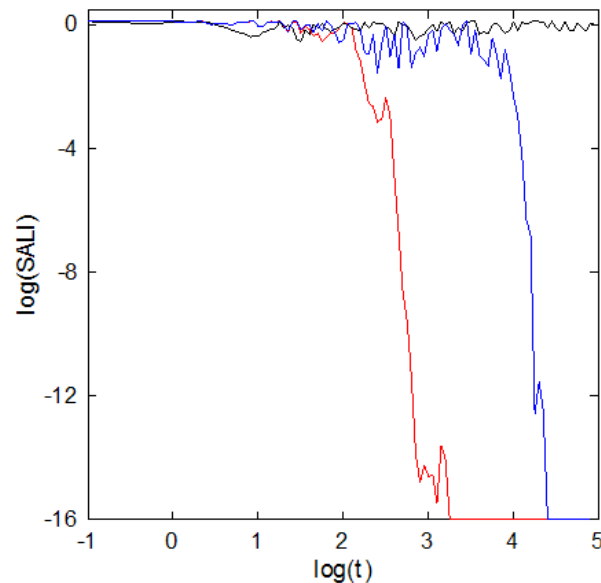
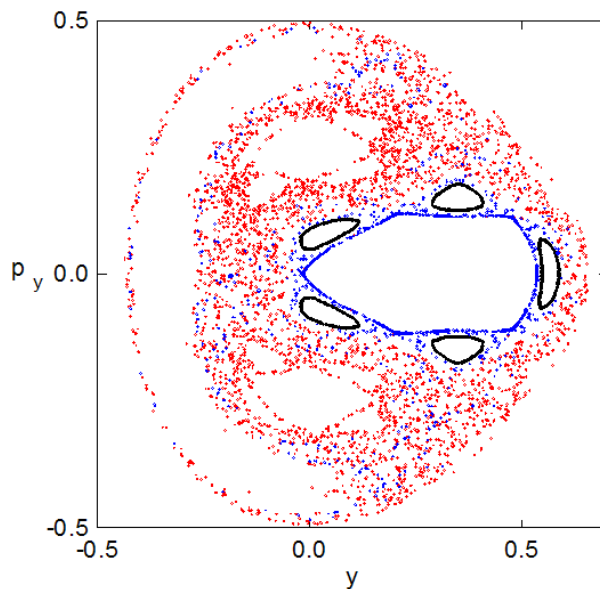
$$H = \frac{1}{2}(p_x^2 + p_y^2) + \frac{1}{2}(x^2 + y^2) + x^2y - \frac{1}{3}y^3$$

For $E=1/8$ we consider the orbits with initial conditions:

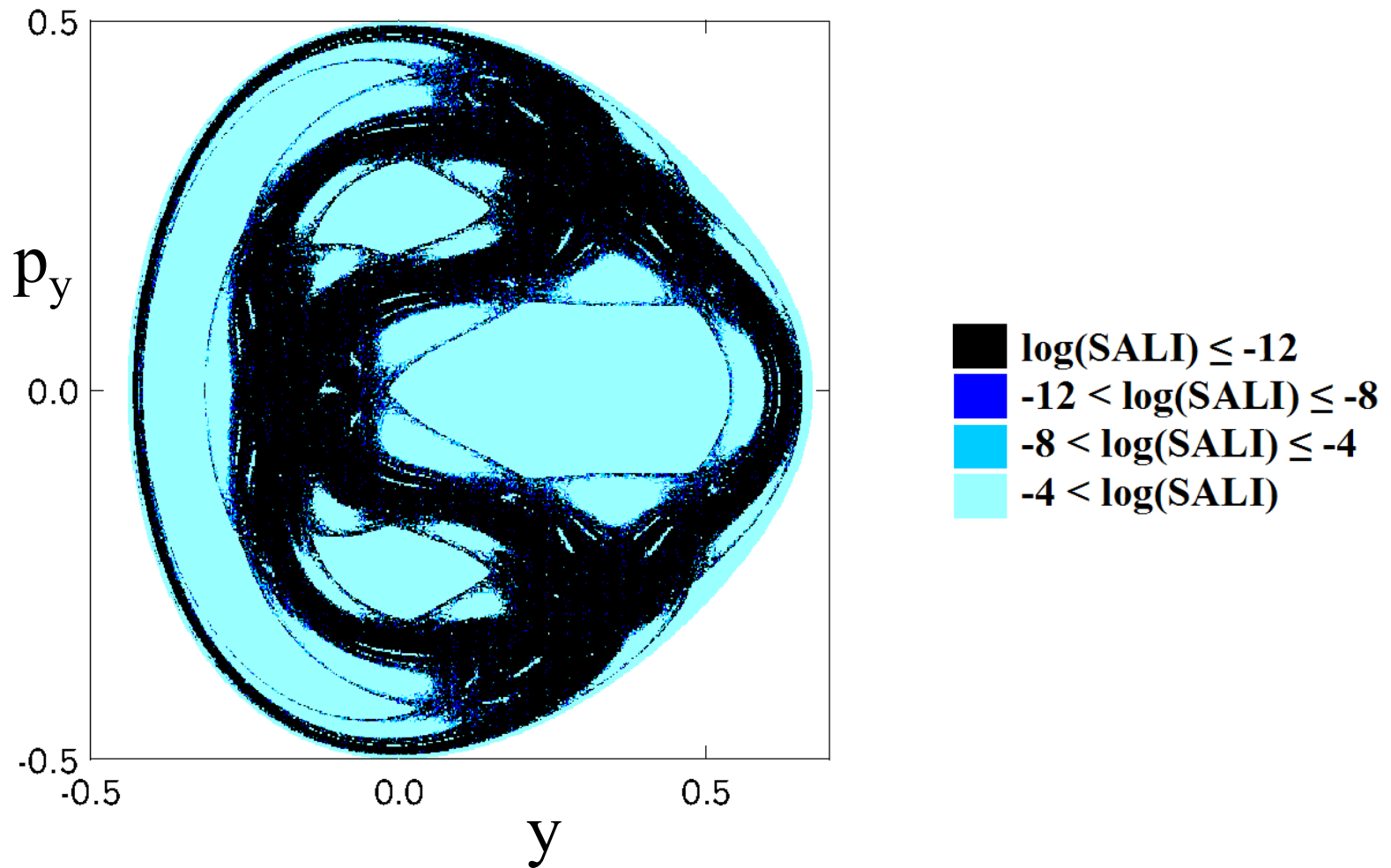
Regular orbit, $x=0, y=0.55, p_x=0.2417, p_y=0$

Chaotic orbit, $x=0, y=-0.016, p_x=0.49974, p_y=0$

Chaotic orbit, $x=0, y=-0.01344, p_x=0.49982, p_y=0$



SALI – Hénon-Heiles system



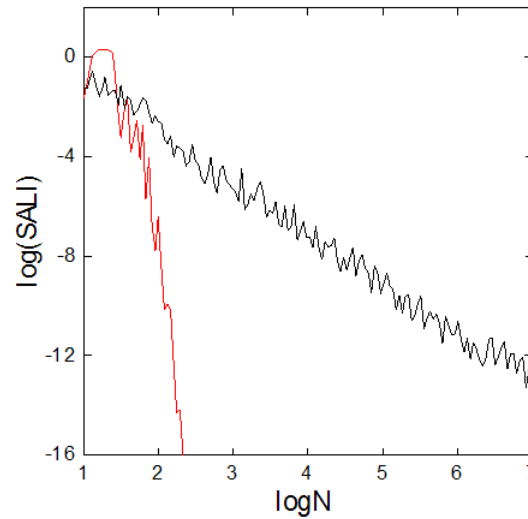
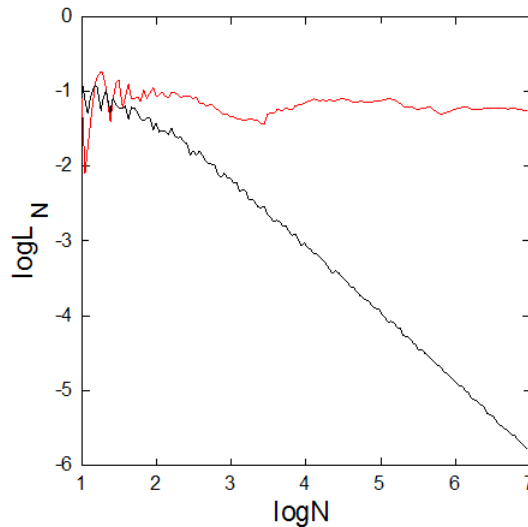
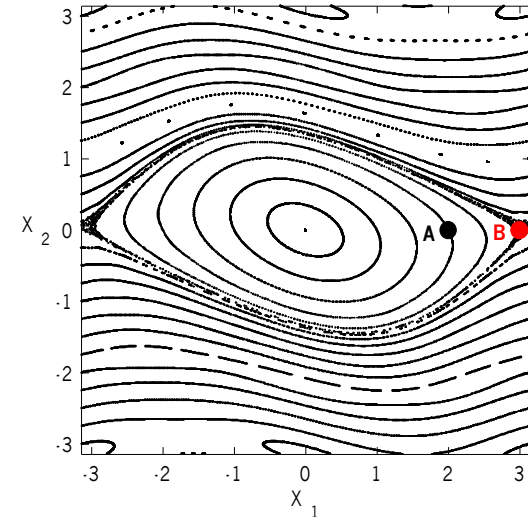
Applications – 2D map

$$\begin{aligned}x'_1 &= x_1 + x_2 \\x'_2 &= x_2 - \nu \sin(x_1 + x_2)\end{aligned}\quad (\text{mod } 2\pi)$$

For $\nu=0.5$ we consider the orbits:

regular orbit A with initial conditions $x_1=2, x_2=0$.

chaotic orbit B with initial conditions $x_1=3, x_2=0$.



Behavior of the SALI

2D maps

SALI $\rightarrow 0$ both for regular and chaotic orbits

following, however, completely different time rates which allows us to distinguish between the two cases.

Hamiltonian flows and multidimensional maps

SALI $\rightarrow 0$ for chaotic orbits

SALI $\rightarrow \text{constant} \neq 0$ for regular orbits

Using LDs to quantify chaos

We consider orbits on a finite grid of an $n(\geq 1)$ -dimensional subspace of the $N(\geq n)$ -dimensional phase space of a dynamical system and their LDs.

Any non-boundary point x in this subspace has $2n$ nearest neighbors

$$y_i^\pm = x \pm \sigma^{(i)} e^{(i)}, \quad i = 1, 2, \dots, n,$$

where $e^{(i)}$ is the i th usual basis vector in \mathbb{R}^n and $\sigma^{(i)}$ is the distance between successive grid points in this direction.

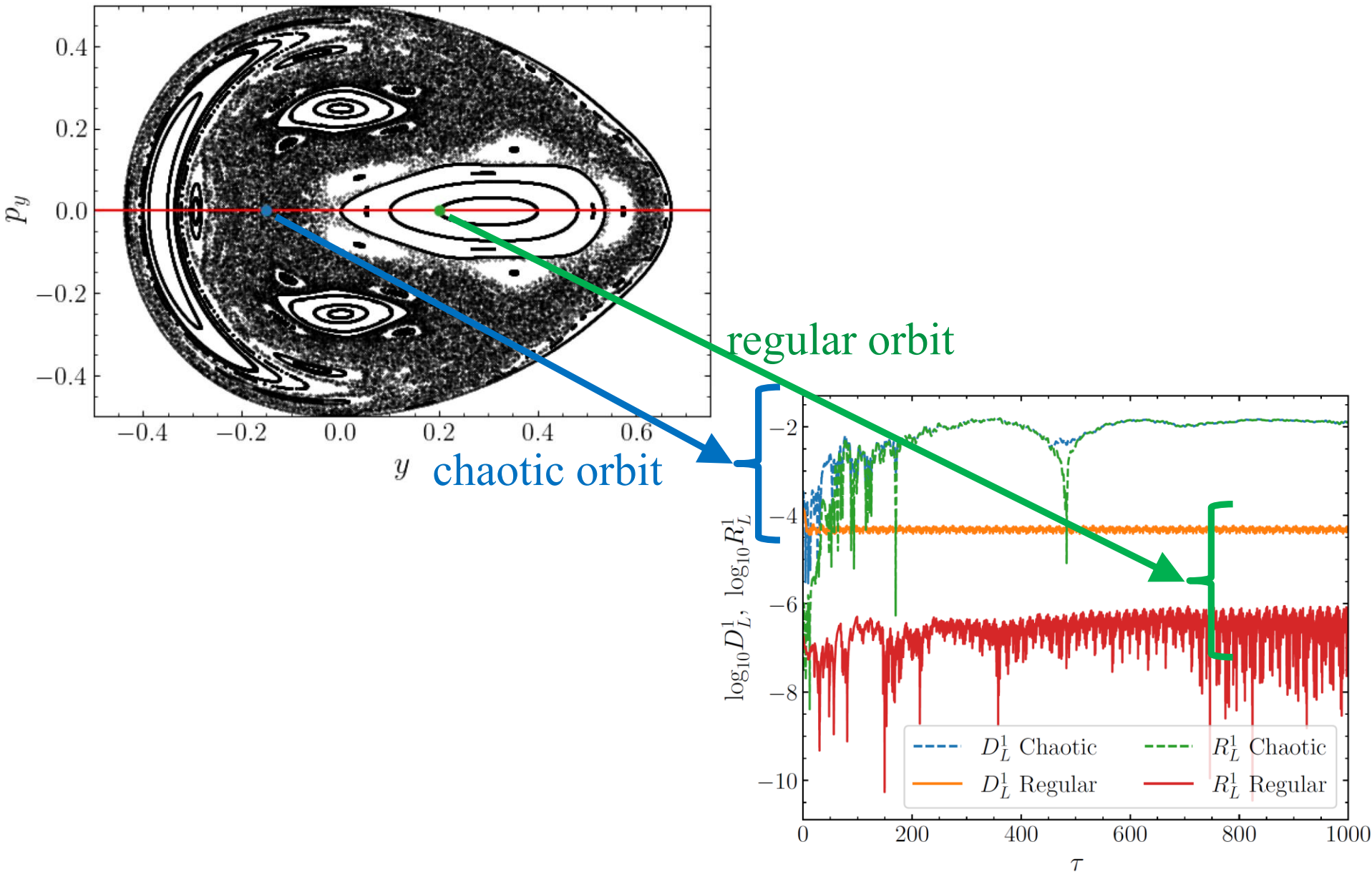
The **difference** D_L^n of neighboring orbits' LDs:

$$D_L^n(x) = \frac{1}{2n} \sum_{i=1}^n \frac{|LD^f(x) - LD^f(y_i^+)| + |LD^f(x) - LD^f(y_i^-)|}{LD^f(x)}.$$

The **ratio** R_L^n of neighboring orbits' LDs:

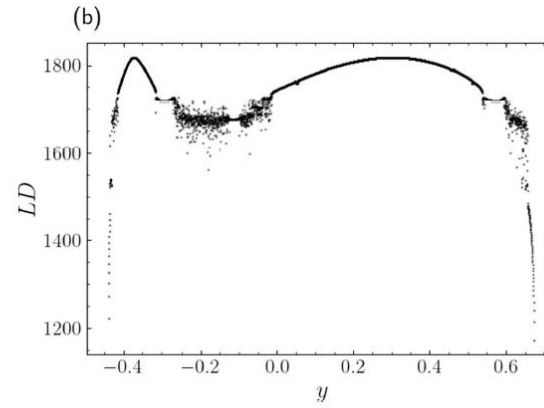
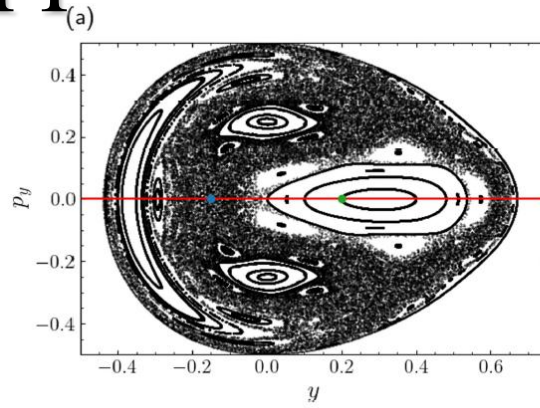
$$R_L^n(x) = \left| 1 - \frac{1}{2n} \sum_{i=1}^n \frac{LD^f(y_i^+) + LD^f(y_i^-)}{LD^f(x)} \right|.$$

Application: Hénon-Heiles system



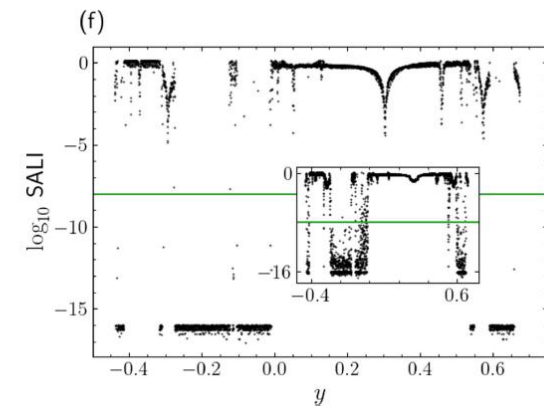
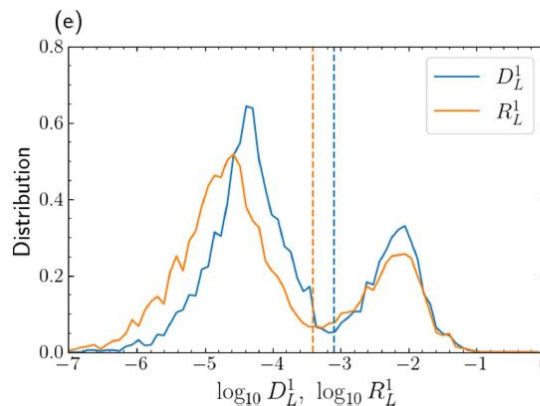
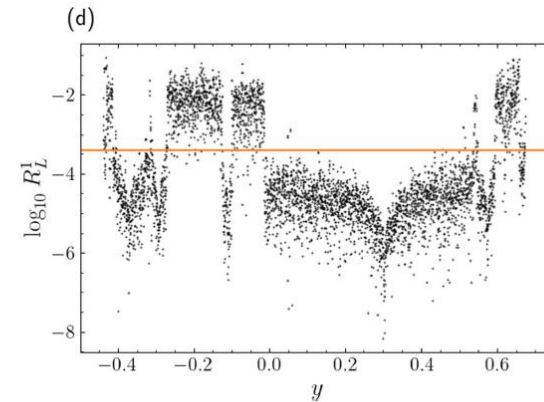
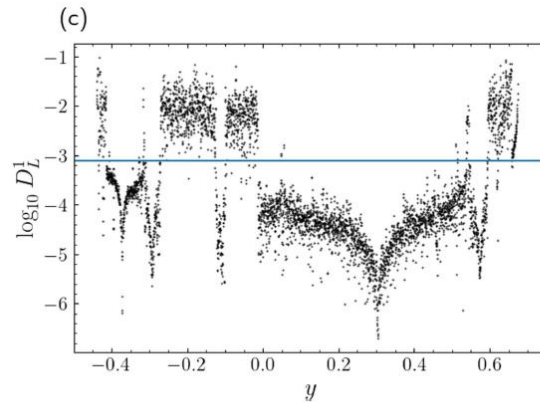
Application: Hénon-Heiles system

$H=1/8$



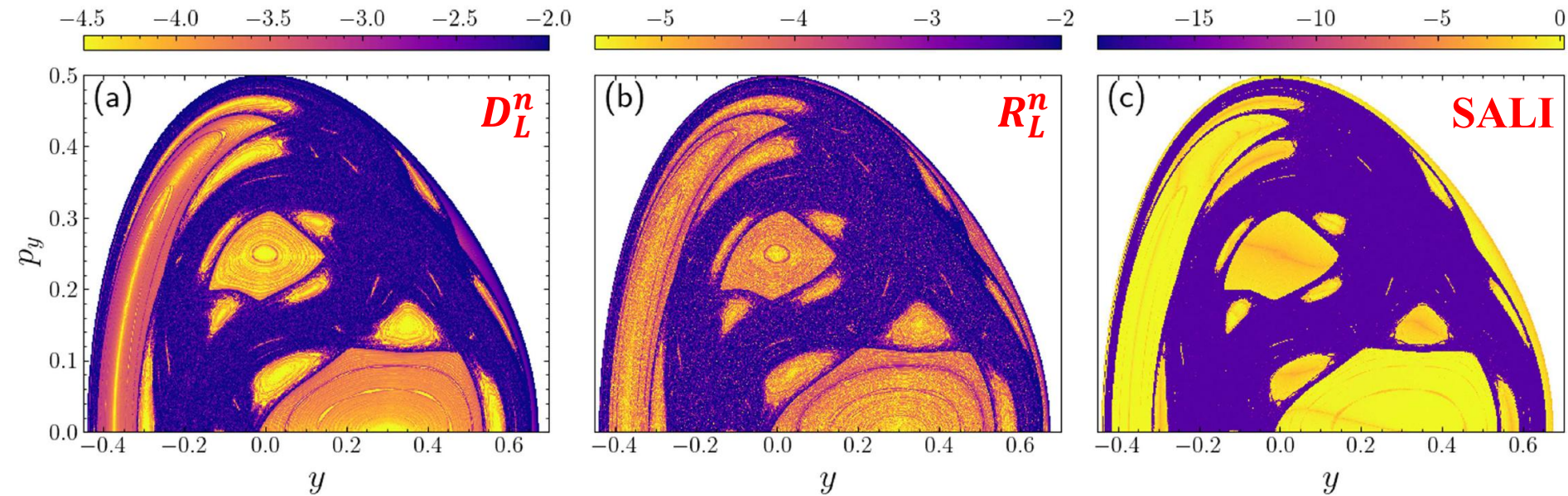
Variation of LDs with regard to initial conditions.
 regular regions: smooth
 chaotic regions: erratic
 [also see Montes et al., Commun. Nonlin. Sci. Num. Simul. (2021)]

LDs for $\tau=10^3$

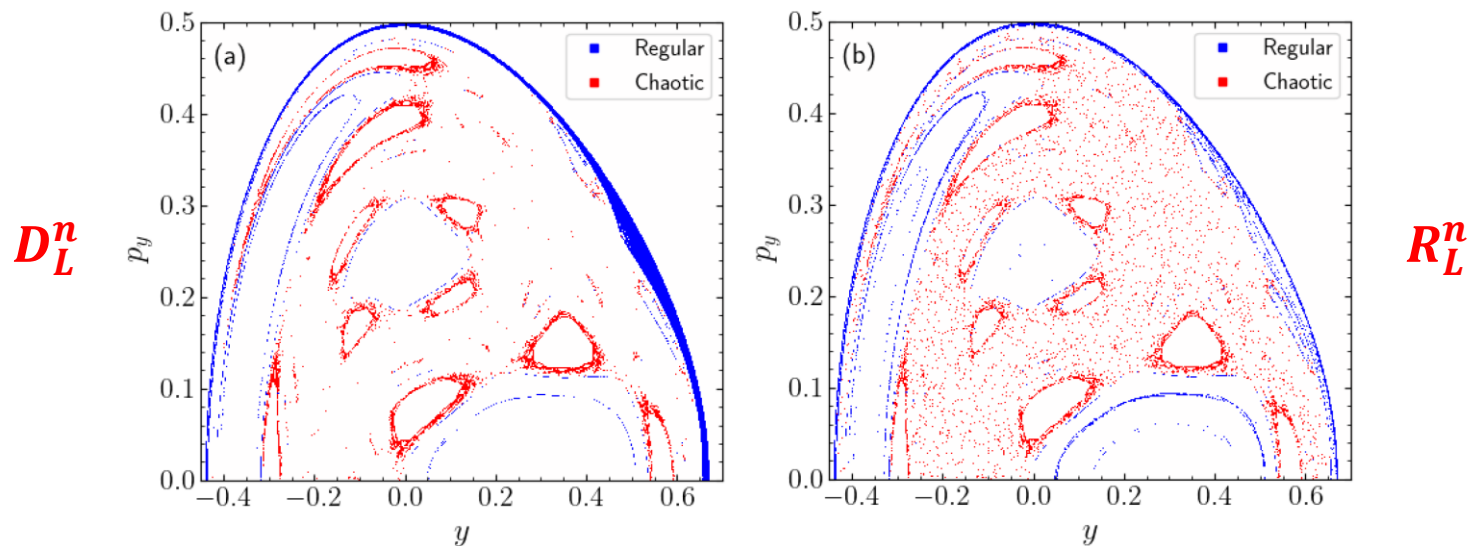


SALI for $\tau=10^6$
 (inset $\tau=10^3$)

Application: Hénon-Heiles system



Misclassified orbits ($< 10\%$)



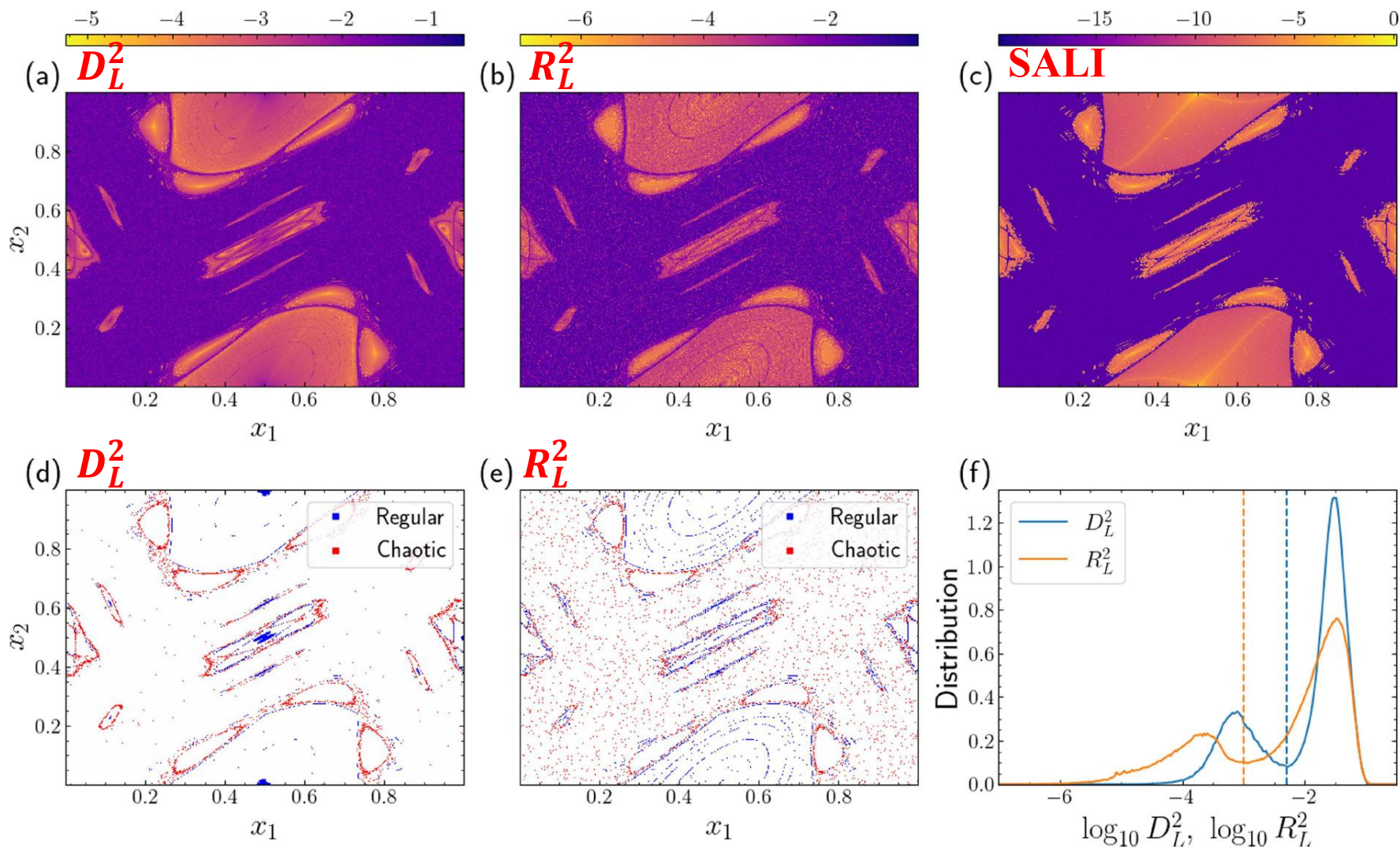
Application: 2D Standard map

$$\begin{aligned} x'_1 &= x_1 + x'_2 \\ x'_2 &= x_2 + \frac{K}{2\pi} \sin(2\pi x_1) \pmod{1} \end{aligned}$$

We set $K = 1.5$

Thresholds: $\log_{10} D_L^2 = -2.3$, $\log_{10} R_L^2 = -3$ ($T = 10^3$)

$\log_{10} \text{SALI} = -12$ ($T = 10^5$)

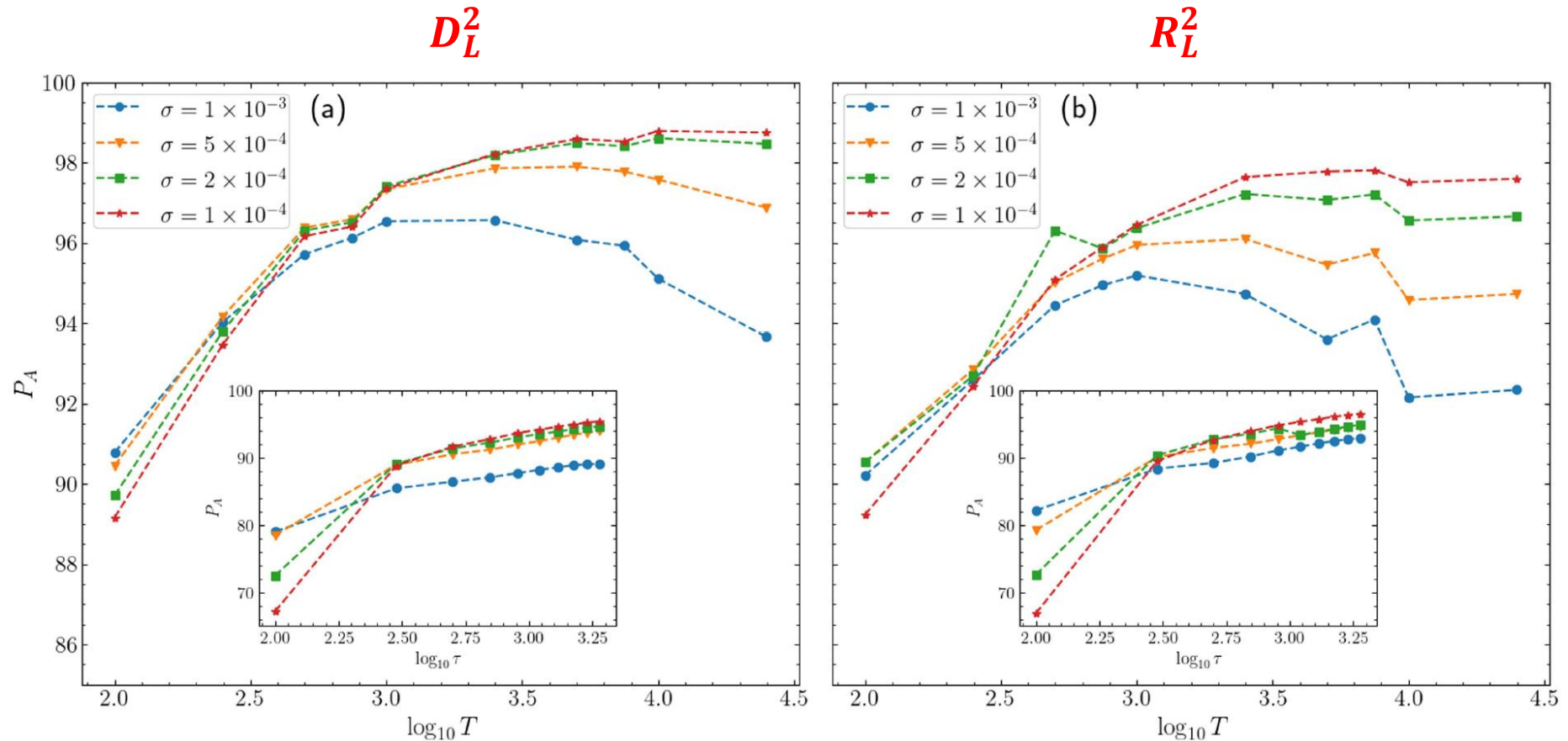


Effect of grid spacing (σ) and final integration time (T, τ)

P_A : percentage of correctly characterized orbits

Main plots: 2D Standard map

Insets: Hénon-Heiles system

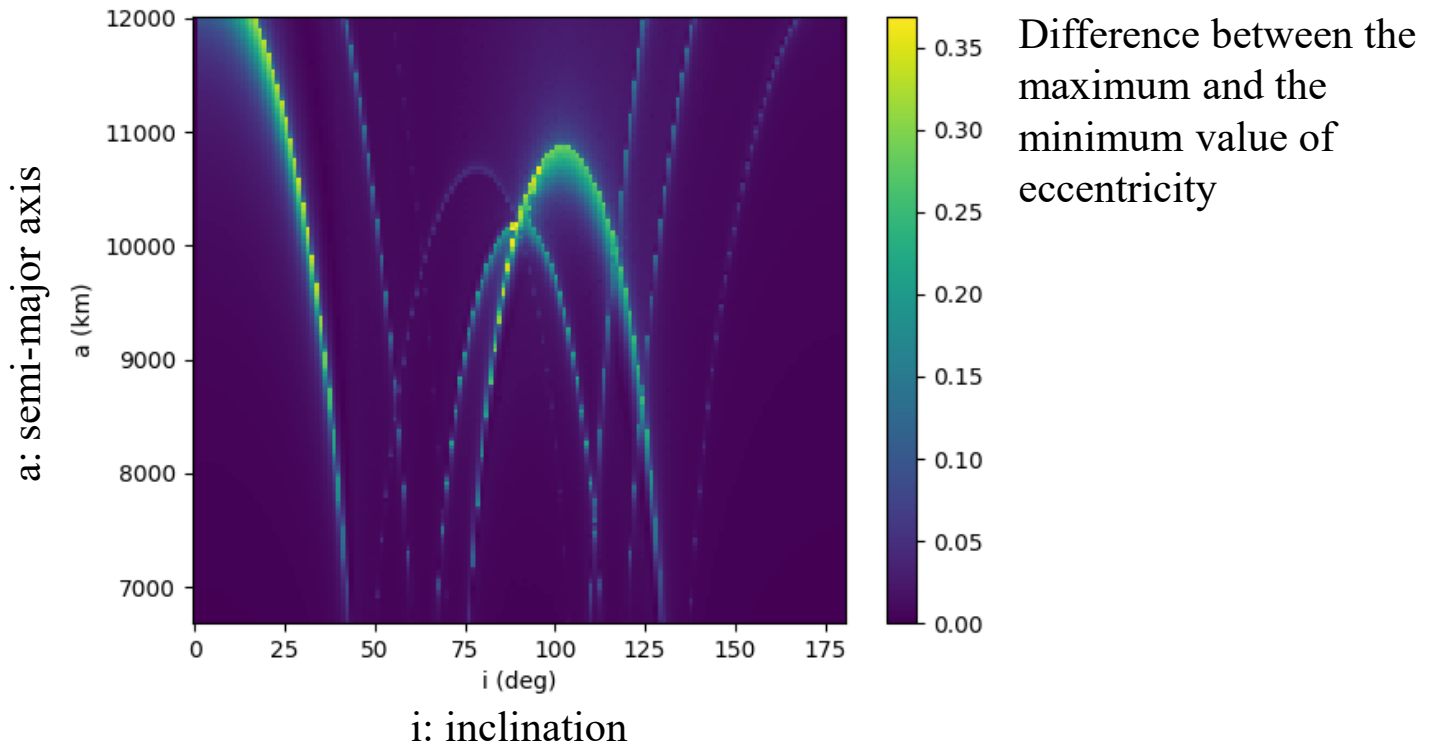


Current work

Investigation of the **dynamics of a satellite around the Earth**.

Following the work of Gkolias et al., Cel. Mech. Dyn. Astron. (2020), we consider a Hamiltonian system consisting of terms describing:

- ✓ **the two body problem** (unperturbed, integrable part), and
- ✓ **perturbation terms** due to
 - **the solar radiation pressure**, and
 - **Earth's oblateness**



Summary

- ✓ We introduced and successfully implemented computationally efficient ways to **effectively identify chaos** in conservative dynamical systems **from the values of LDs at neighboring initial conditions**.
- ✓ From the distributions of the indices' values we determine appropriate **threshold values**, which allow the characterization of orbits as regular or chaotic.
- ✓ All indices **faced problems** in correctly revealing the nature of some orbits mainly **at the borders of stability islands**.
- ✓ All indices show **overall very good performance**, as their classifications are in accordance with the ones obtained by **the SALI (which is a very efficient and accurate chaos indicator)** at a level of at least 90% agreement.
- ✓ **Advantages:**
 - **Easy to compute** (actually only the forward LDs are needed).
 - **No need to know and to integrate the variational equations.**

References

- Madrid & Mancho, Chaos, 19, 013111 (2009)
- Mendoza & Mancho, PRL, 105, 038501 (2010)
- Mancho, Wiggins, Curbelo & Mendoza, Commun. Nonlin. Sci. Num. Simul., 18, 3530 (2013)
- Lopesino, Balibrea, Wiggins & Mancho, Commun. Nonlin. Sci. Num. Simul., 27, 40 (2015)
- Lopesino, Balibrea-Iniesta, García-Garrido, Wiggins & Mancho, Int. J. Bifurc. Chaos, 27, 1730001 (2017)
- Agaoglou, Aguilar-Sanjuan, García-Garrido, González-Montoya, Katsanikas, Krajňák, Naik & Wiggins, ‘Lagrangian descriptors: Discovery and quantification of phase space structure and transport’, <https://doi.org/10.5281/zenodo.3958985> (2020)
- Montes, Revuelta & Borondo, Commun. Nonlin. Sci. Num. Simul., 102, 105860 (2021)
- Daquin, Pédenon-Orlanducci, Agaoglou, García-Sánchez & Mancho, Physica D, 442, 133520 (2022)
- S., J. Phys. A, 34, 10029 (2001)
- S., Antonopoulos, Bountis & Vrahatis, J. Phys. A, 37, 6269 (2004)
- S. & Manos, Lect. Notes Phys., 915, 129 (2016)

Hillebrand, Zimmer, Ngapasare, Katsanikas, Wiggins & S.: Chaos, 32, 123122 (2022),
‘Quantifying chaos using Lagrangian descriptors’

Zimmer, Ngapasare, Hillebrand, Katsanikas, Wiggins & S.: Physica D, 453, 133833 (2023),
‘Performance of chaos diagnostics based on Lagrangian descriptors. Application to the 4D standard map’